

Predictions of slow longitudinal mode 1 instability in storage rings with harmonic cavities

Murilo B. Alves
Brazilian Synchrotron Light Laboratory
LNLS-SIRIUS



Outline

1. Introduction
2. Prediction with macroparticle tracking
3. Prediction based on instability theory
4. Conclusions

Mode 1 Instability

- 4th generation storage rings with passive harmonic cavities (HCs) for bunch lengthening
- Longitudinal plane. Uniform filling pattern. NC and SC passive HCs
- Coherent oscillation of coupled-bunch mode 1 with low frequency (few Hz)
- Driven above a critical harmonic voltage, may limit performance with HC
- Predicted in simulations during HC studies (ALS-U¹, HALF², SOLEIL-II³, DIAMOND-II⁴). Observed experimentally at MAX-IV⁵

1) M. Venturini, PRAB **21** 114404 (2018)

2) T. He, et al. PRAB **25** 024401 (2022)

3) A. Gamelin, Harmonic cavity studies for the SOLEIL Upgrade, I.FAST Workshop 2022

4) T. Olsson, Collective Effects in the Diamond-II Storage Ring. LEL Workshop 2022

5) F. Cullinan, et al, Longitudinal Beam Dynamics in Ultra-low Emittance Rings, I.FAST Workshop 2022



CNPq

Recent developments to predict the mode 1 instability

2nd HarmonLIP, Grenoble, 18-20 March 2024

MINISTRY OF
SCIENCE TECHNOLOGY
AND INNOVATION



M. Venturini, 2018



CNPEM

- Ring: ALS-U. Normal-conducting cavities
- Perturbation theory (mode analysis) in **arbitrary longitudinal potential for any coupled-bunch mode.** Elements of the theory in S. Krinsky, 1985; A. Mosnier, 1999; Y. Cai 2011
- Focused on dipolar motion of mode 0 (Robinson) and mode 1. Simplified formulas for quadratic (single-rf) and quartic potential (flat-potential)
- Predicted **low coherent oscillation frequency** for mode 1
- Highlighted mode 1 instability **driven by the imaginary part of HCs impedance**
- Mode 1 ruled out the option of reusing ALS HCs

PHYSICAL REVIEW ACCELERATORS AND BEAMS **21**, 114404 (2018)

Passive higher-harmonic rf cavities with general settings and multibunch instabilities in electron storage rings

Marco Venturini*

Lawrence Berkeley National Laboratory, University of California, Berkeley, California 94720, USA



(Received 19 September 2018; published 27 November 2018)


- Ring: HALF. Superconducting cavities
- **Periodic transient beam loading (PTBL)** predicted with **macroparticle tracking** and **semi-analytical equilibrium solver**. Explored the characteristics of PTBL
- GPU accelerated tracking with thousands of macroparticles for millions of turns
- Thresholds from non-convergence of semi-analytical equilibrium agrees with tracking. Possible due to the property of **low coherent frequency of PTBL**
- Steady-state time-domain perturbation approach
- Applied a **mode 1 phase perturbation** to the system and studied if it is amplified
- **Analytical formula** with results that agree well with macroparticle tracking

PHYSICAL REVIEW ACCELERATORS AND BEAMS **25**, 024401 (2022)

Periodic transient beam loading effect with passive harmonic cavities in electron storage rings

Tianlong He^{✉*}, Weiwei Li,[†] Zhenghe Bai,[‡] and Lin Wang

National Synchrotron Radiation Laboratory, University of Science and Technology of China, Hefei, Anhui, 230029, China

 (Received 13 October 2021; accepted 18 January 2022; published 3 February 2022)

PHYSICAL REVIEW ACCELERATORS AND BEAMS **25**, 094402 (2022)

Novel perturbation method for judging the stability of the equilibrium solution in the presence of passive harmonic cavities

Tianlong He^{✉*}

National Synchrotron Radiation Laboratory, University of Science and Technology of China, Hefei, Anhui, 230029, People's Republic of China

 (Received 15 June 2022; accepted 15 September 2022; published 29 September 2022)

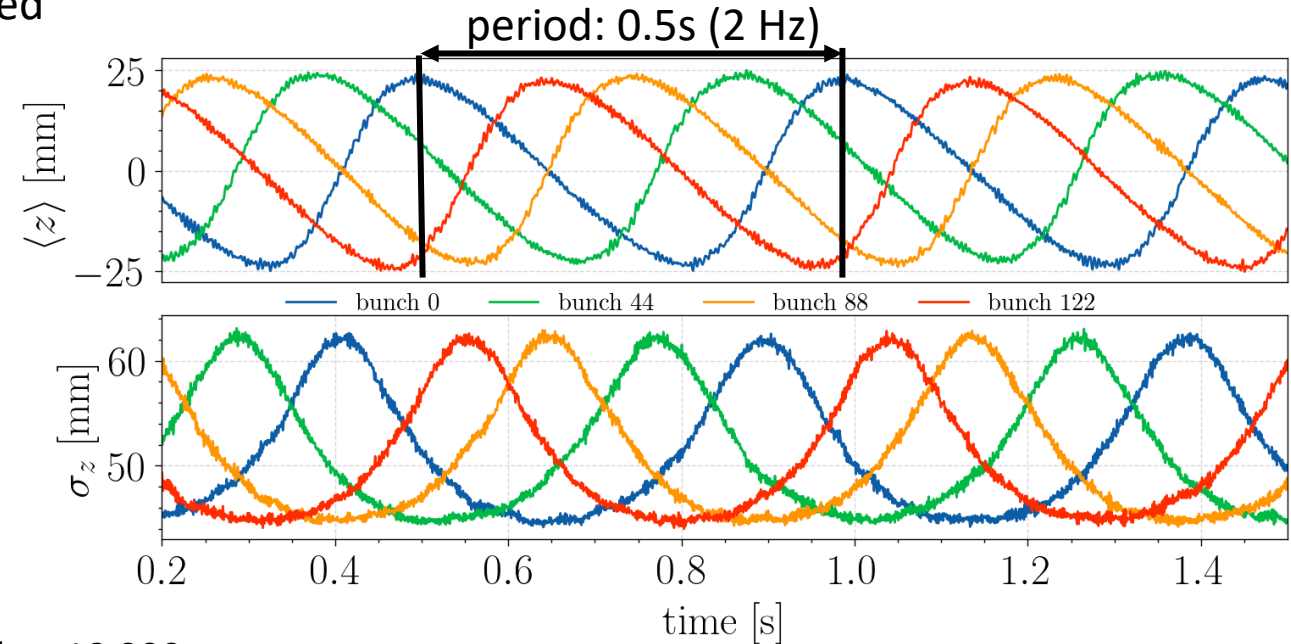
Prediction with macroparticle tracking

Validation of tracking code

- Longitudinal tracking code in python3 developed at LNLS*
- Test a case where mode 1 instability is already expected

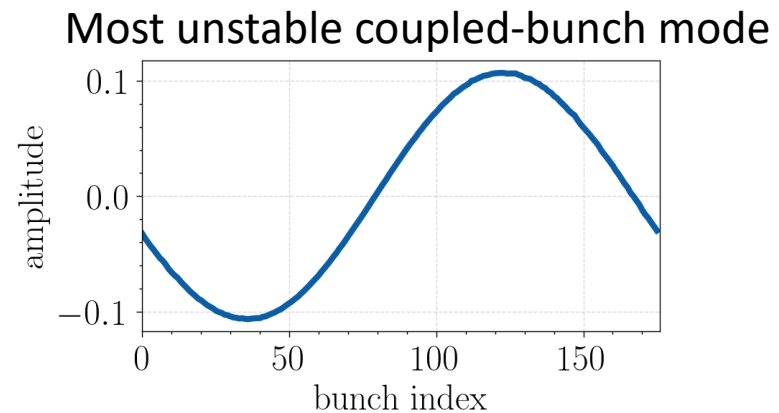
MAX-IV Parameters

Energy	E_0	3 GeV
Current	I_0	300 mA
RF Frequency	f_{rf}	99.931 MHz
Harmonic Number	h	176
Momentum Compaction	α	3.06×10^{-4}
Energy Spread	σ_δ	7.69×10^{-4}
Bunch Length	σ_z	10.1 mm
Energy Loss	U_0	363.8 keV
Gap Voltage	V_{rf}	1.397 MV
Synchrotron Tune	ν_z	0.001 638
Synchrotron Frequency	f_z	930 Hz
Longitudinal Damping	τ_δ	25.194 ms
HC Type		Passive NC
HC RF harmonic	q	3
HC Shunt Impedance	R_s	8.25 M Ω
HC Quality Factor	Q	20 800
HC R/Q	R/Q	396 Ω



Nr. of Macroparticles: 10 000
Nr. of Turns: 1 500 000

No beam loading from main cavity
Passive HC at flat potential (300mA, RF voltage 1.397 MV)
Detune: 75 kHz
Harmonic voltage: 448 kV

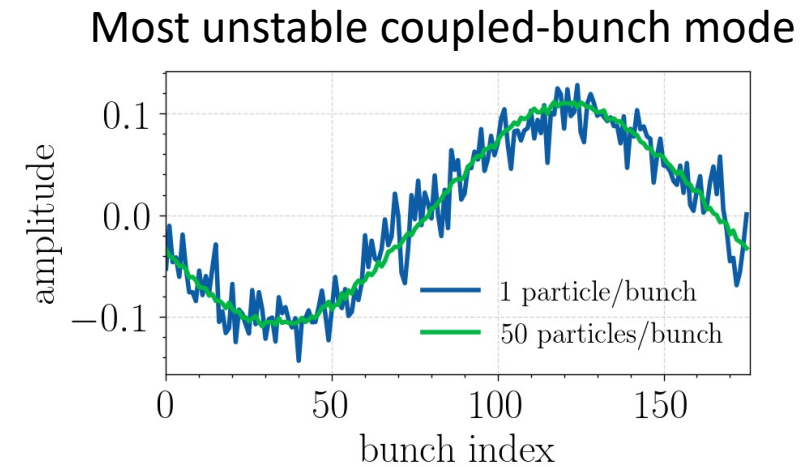
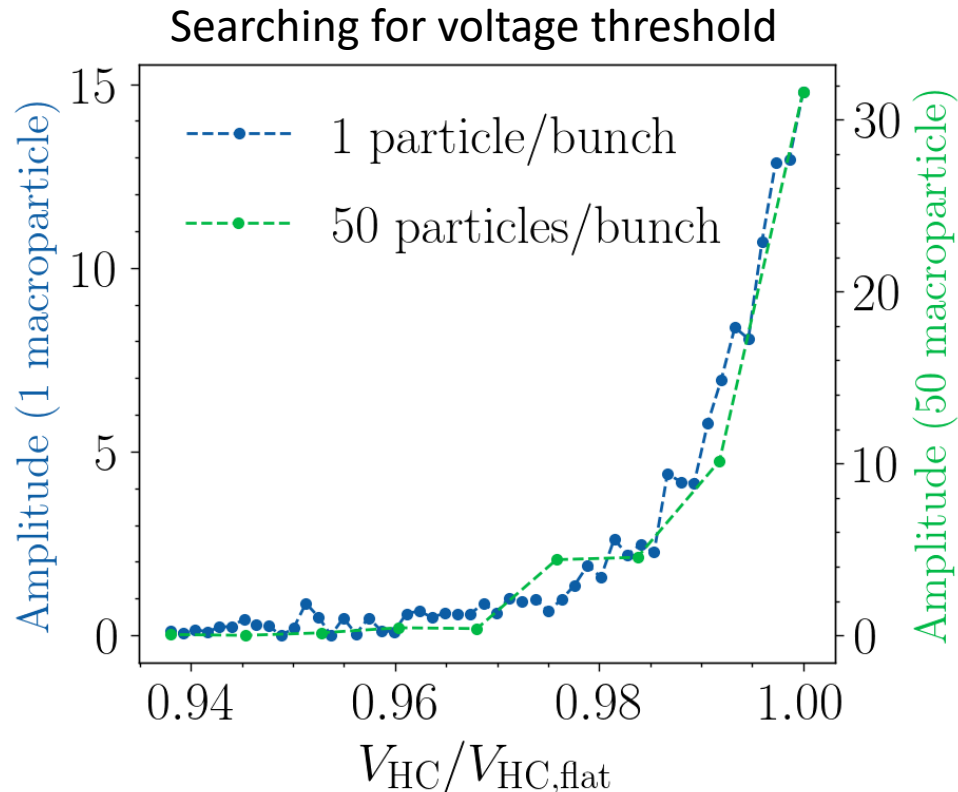


* Open access code

https://github.com/lnls-fac/collective_effects/blob/master/pycolleff/pycolleff/longitudinal_tracking.py

Simplifying the simulation

- Does the instability depend on bunch details? → Tests with reduced number of macroparticles
- Mode 1 instability happens even in the limiting case of 1 macroparticle per bunch!
 - Bunch centroids (point bunch approximation) should capture the underlying mechanism for instability
 - Non-linearities inside bunch should be negligible





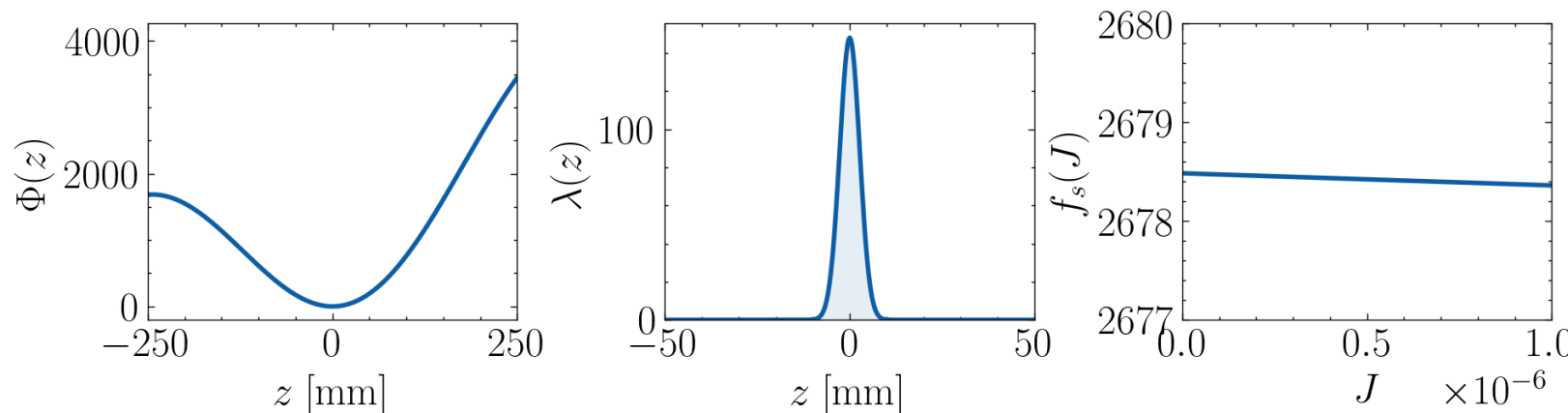
CNPq

Prediction based on instability theory

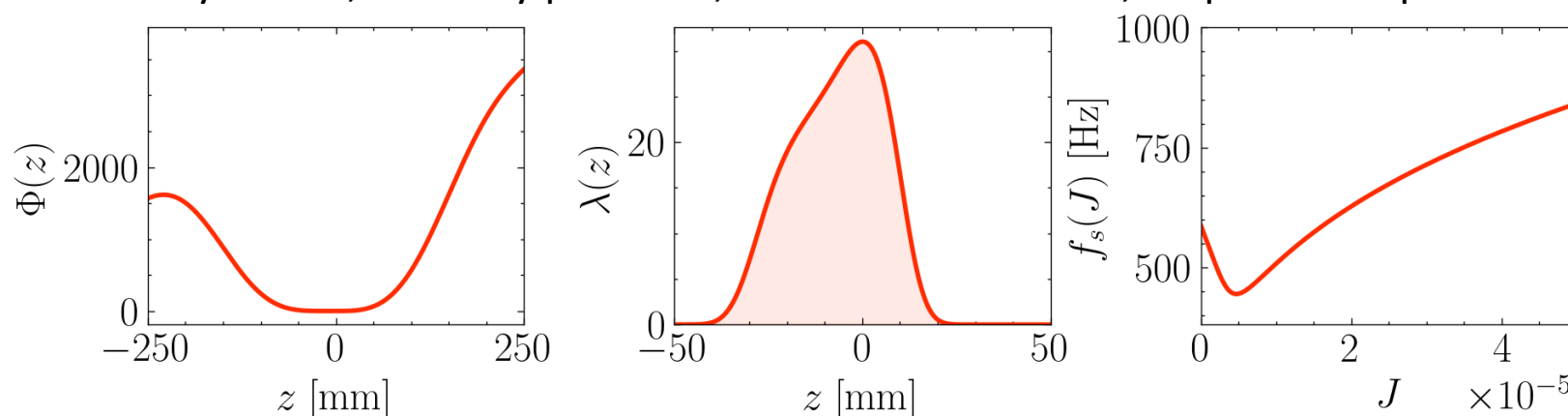
Step 1: solve the equilibrium with HCs

- Self-consistent solution of Haissinski equation
- Equilibrium longitudinal potential and bunch profiles

Single-rf: linear dynamics, quadratic potential, Gaussian bunch, constant synchrotron frequency



Double-rf: non-linear dynamics, arbitrary potential, non-Gaussian bunch, amplitude-dependent synchrotron frequency



Open access code

https://github.com/lnls-fac/collective_effects/blob/master/pycolleff/pycolleff/longitudinal_equilibrium.py

Step 2: evaluate the stability of mode 1



- HC fields introduce non-linearity and modify the bunch profile. Out of scope of linear instability theory for Gaussian bunches. Linear theory can still be useful though
- Stability diagram analysis:
 - Calculate the complex coherent frequencies for the linear system
 - Evaluate the location of unstable modes on the stability boundary (effect of non-linearities)
 - Applied for instabilities with HCs: M. Venturini, PRAB **21** 114404 (2018) and R. Lindberg, PRAB **21** 124402 (2018)
- Hypothesis: the coherent frequency of mode 1 is too low to have enough overlap with spread of incoherent frequencies → Non-linearities/Landau damping effects can be neglected

Goal: apply a linear theory for Gaussian bunches to predict the mode 1 instability

Longitudinal Mode-Coupling Instability (LMCI)

Suzuki, T., Chin, Y., & Satoh, K. (1983). Mode-coupling theory and bunch lengthening in SPEAR II. Particle Accelerators, 13(3-4), 179-198.

Suzuki, T. (1983). Theory of longitudinal bunched-beam instabilities based on the Fokker-Planck equation. Particle Accelerators, 14, 91-108

- Linear single-particle dynamics in polar coordinates (r, ϕ)
Quadratic potential, Gaussian bunch

- Small perturbation from equilibrium drives dynamic wakefields with coherent frequency

$$\psi(r, \phi, t) = \psi_0(r) + \psi_1(r, \phi)e^{-i\Omega t}$$

- Expansion in **azimuthal modes m** $\psi_1(r, \phi) = \sum_{m=-\infty}^{+\infty} R_m(r)e^{im\phi}$
- Expansion in orthogonal polynomials (Laguerre) \rightarrow **radial modes k** associated with each m

- Insert this perturbation to linearize Fokker-Planck equation and solve for coherent frequencies

- Coherent frequencies are eigenvalues of interaction matrix including **mode coupling of azimuthal and radial modes**. Instability if $\text{Im}\Omega > 0$

- Uniform filling. For analysis of coupled-bunch mode ℓ sample impedance at $(p + \ell)\omega_0$. For mode 1, $\ell = 1$

Open access code

https://github.com/lnls-fac/collective_effects/blob/master/pycolleff/pycolleff/colleff.py

Adaptations for the analysis

- Calculation in a fictitious quadratic potential with longer bunch

$$\sigma_{z,0} \rightarrow \sigma_{z,HC}$$

- Incoherent average synchrotron frequency is also an input

$$\text{Compatible with bunch length } \omega_s \sigma_z = \alpha C \sigma_\delta$$

Local frequency averaged by z-distribution

$$\langle \omega_s \rangle_z = \int_{-\infty}^{\infty} \lambda(z) \sqrt{-\frac{\alpha hc}{2\pi E_0 \omega_{rf}} V'(z)} dz$$

Global frequency averaged by action-distribution

$$\langle \omega_s \rangle_J = 2\pi \int_0^{\infty} \Psi(J) \omega_s(J) dJ$$

Typically lower values

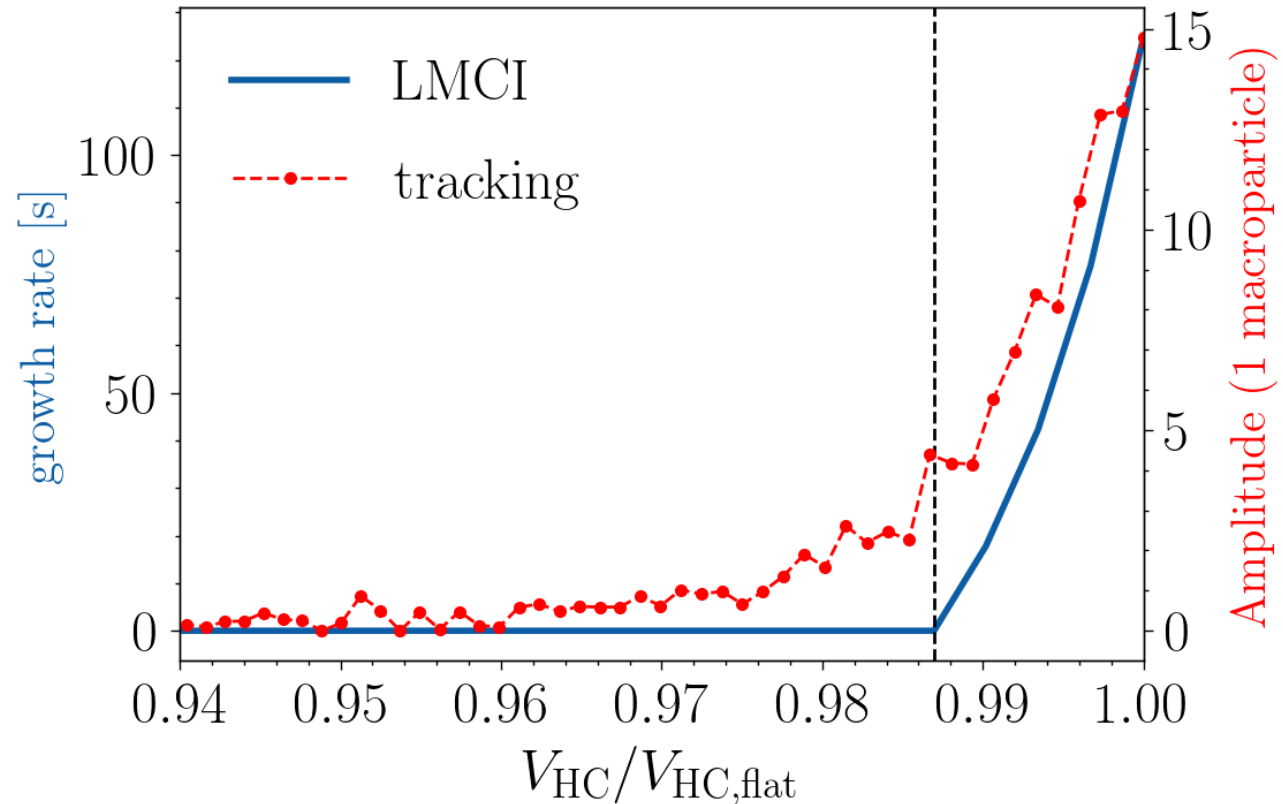
Similar values

- All the following results were computed with HC-modified bunch length + compatible incoherent synchrotron frequency

LMCI results

LMCI vs. Tracking for MAX-IV parameters

- 300mA
- Main RF voltage: 1.397 MV
- 3 HCs
- Agrees with tracking → mode 1 unstable at flat potential

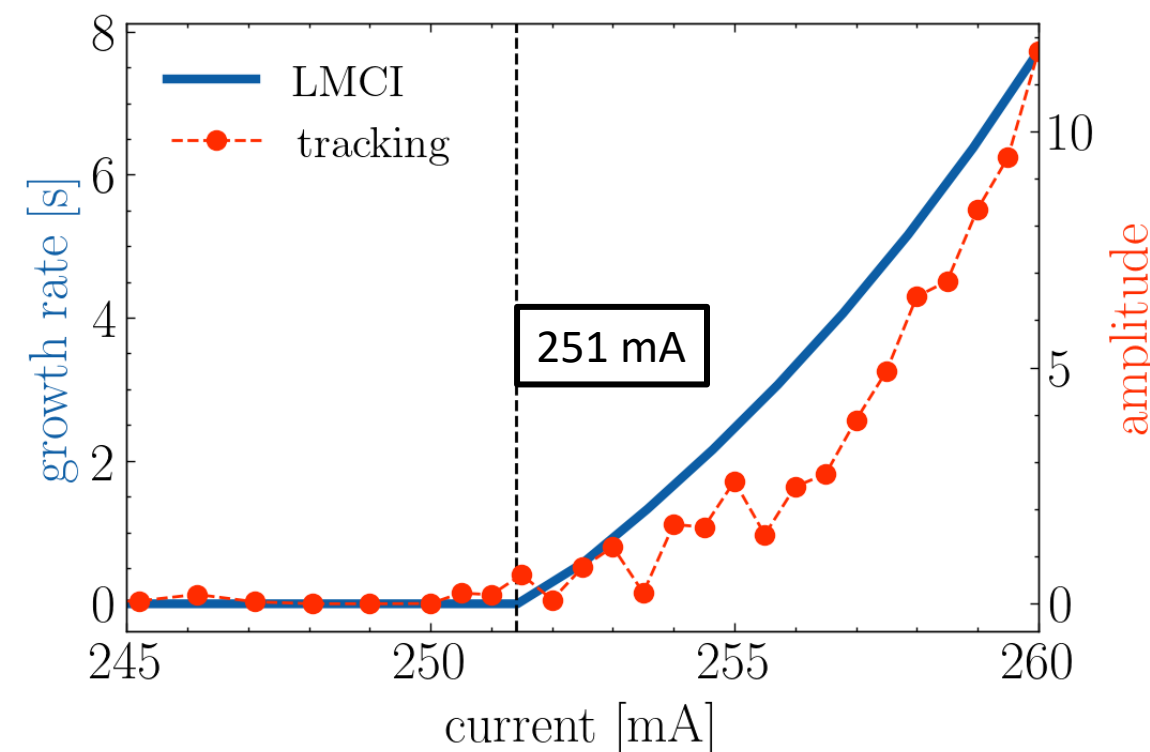


LMCI vs. Tracking for HALF parameters

Parameters for HALF from [1]

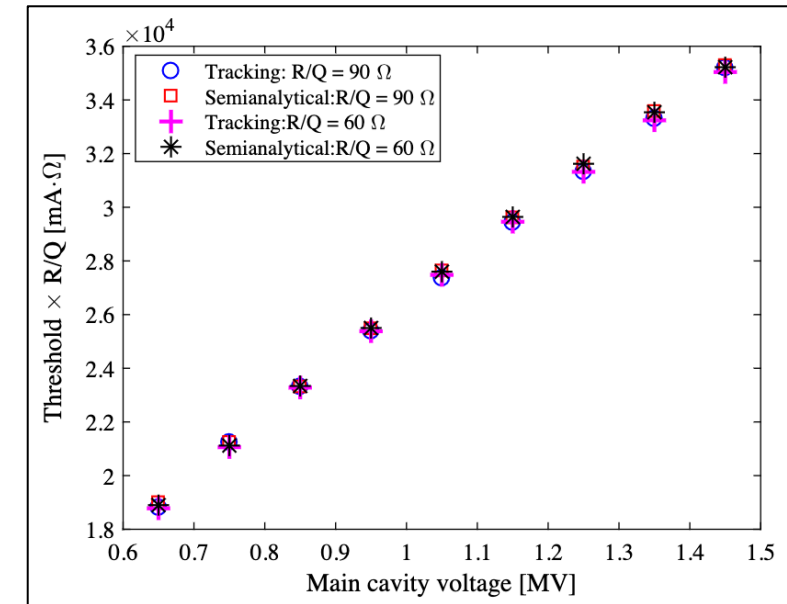
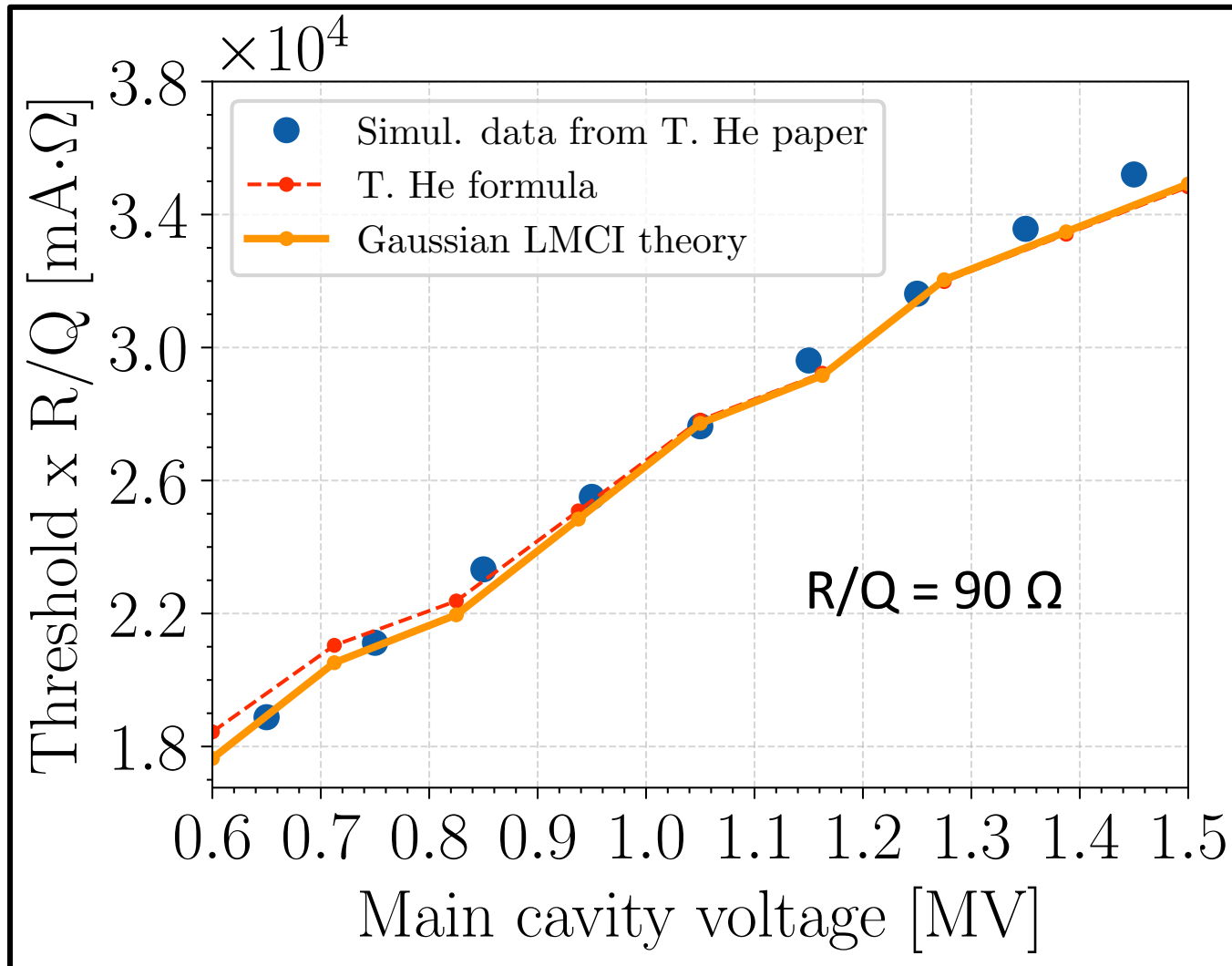
Parameter	Value
Ring circumference	480 m
Beam energy	2.2 GeV
Average beam current	350 mA
Longitudinal damping time	22.7 ms
Momentum compaction factor	8.1×10^{-5}
Natural energy spread	6.45×10^{-4}
Natural bunch length	6.76 ps
Harmonic number	800
Energy loss per turn	198.8 keV
Voltage of MC	0.85 MV
R/Q of 3rd-HC	90 Ω
Quality factor of 3rd-HC	5×10^5

[1] T. He, et al. PRAB **25** 024401 (2022)



259 mA is the reported threshold in [1]
3% difference between the results

LMCI vs. Tracking/Semi-analytical for HALF

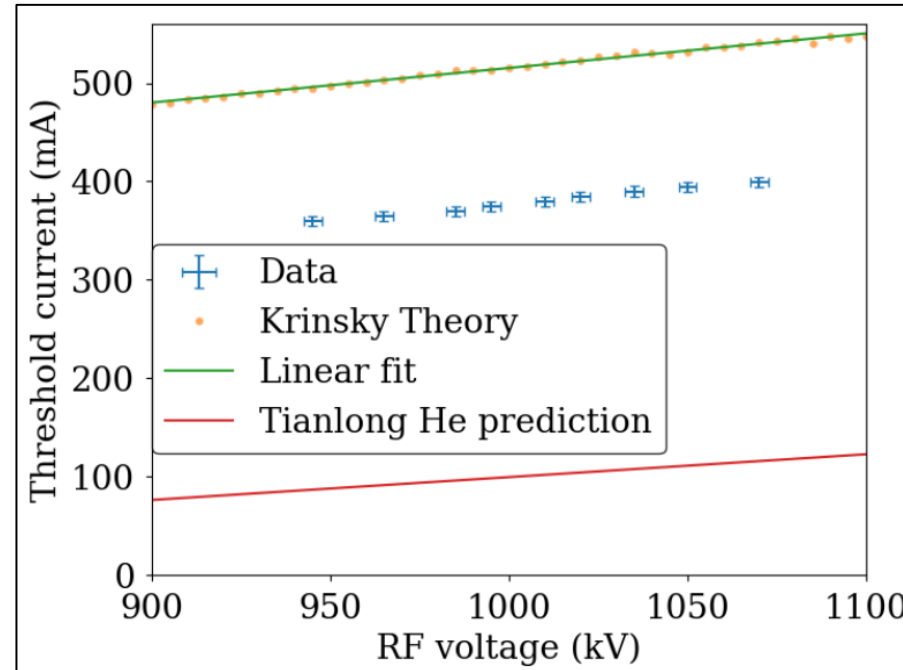


Simulated data from
 T. He, et al. PRAB **25** 024401 (2022)
 Results from STABLE macroparticle
 tracking code and semi-analytical
 method

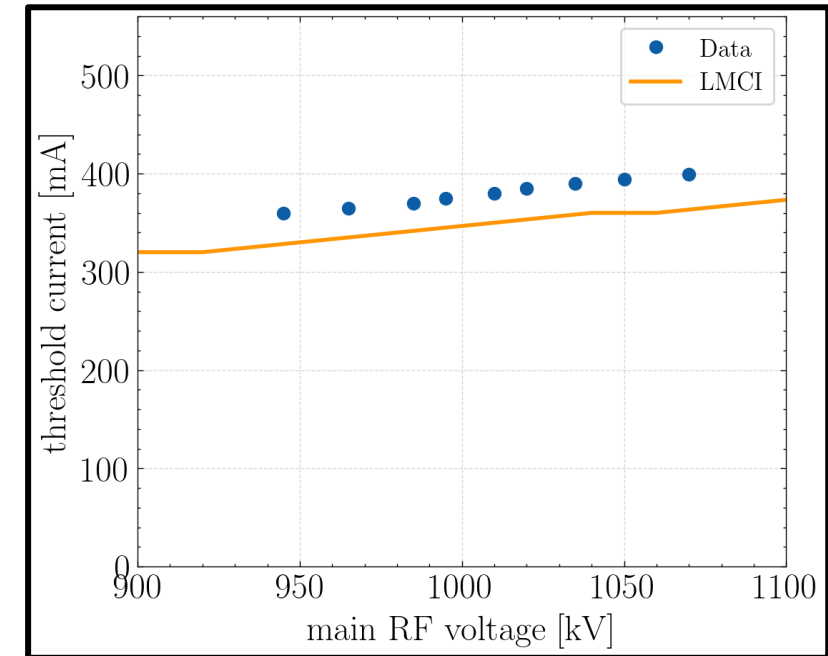
Formula (Eq. 24) from
 T. He, PRAB **25** 094402 (2022)

LMCI vs. Experimental data from MAX-IV

- 2 HCs tuned to flat potential voltage

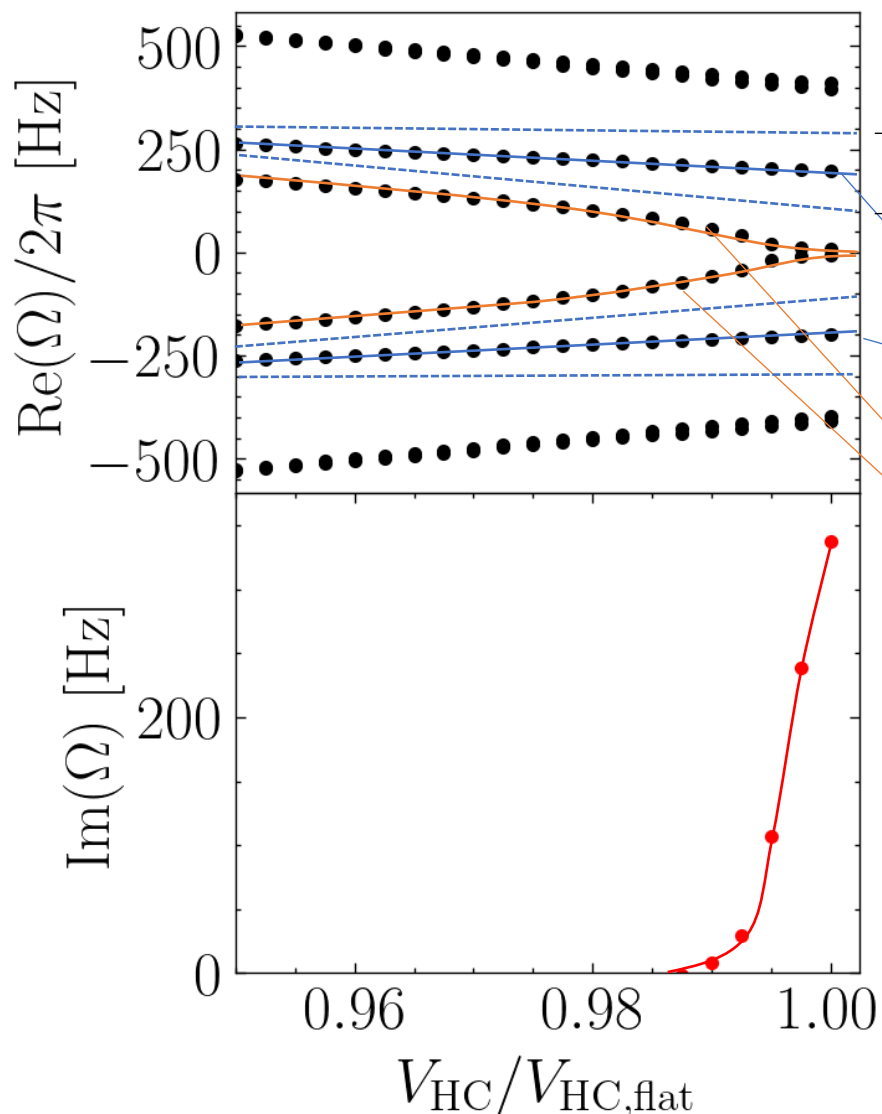


F. Cullinan, et al - Mode 1 Instabilities in Fourth Generation Storage Rings - I.FAST 9th Low Emittance Rings Workshop (2024)



Linear LMCI result agrees with measured values with difference < 10%

Instability mechanism



The hypothesis of non-overlapping coherent and incoherent frequencies is verified

Spread of incoherent synchrotron frequency

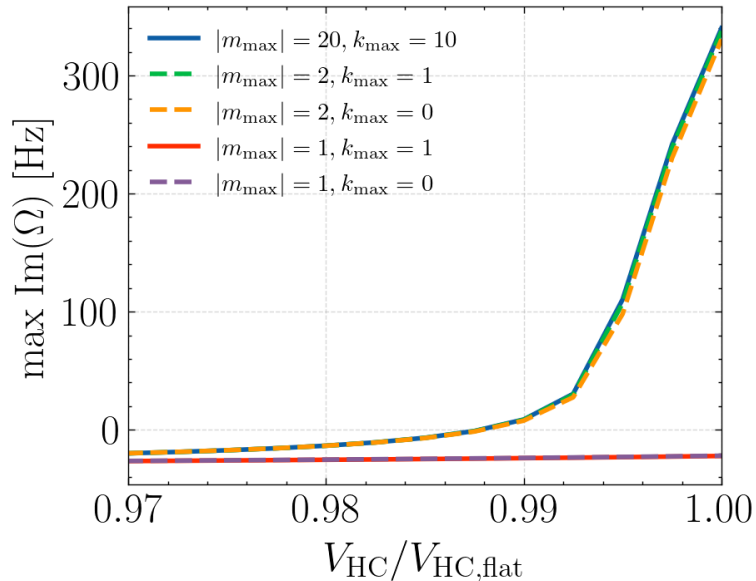
Radial modes following the incoherent frequency $\pm\omega_s$

Radial modes with additional shift $\omega_s - \Delta\Omega$ induced by the imaginary part of the HC impedance

Low dipole coherent frequency for mode 1, very weak longitudinal focusing, instability builds up

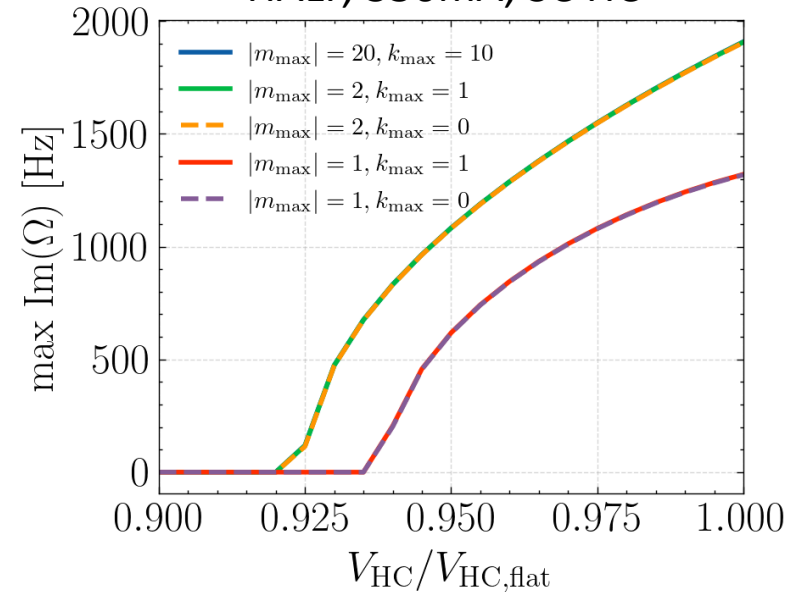
Tests of convergence

MAX-IV, 400mA, NC HC



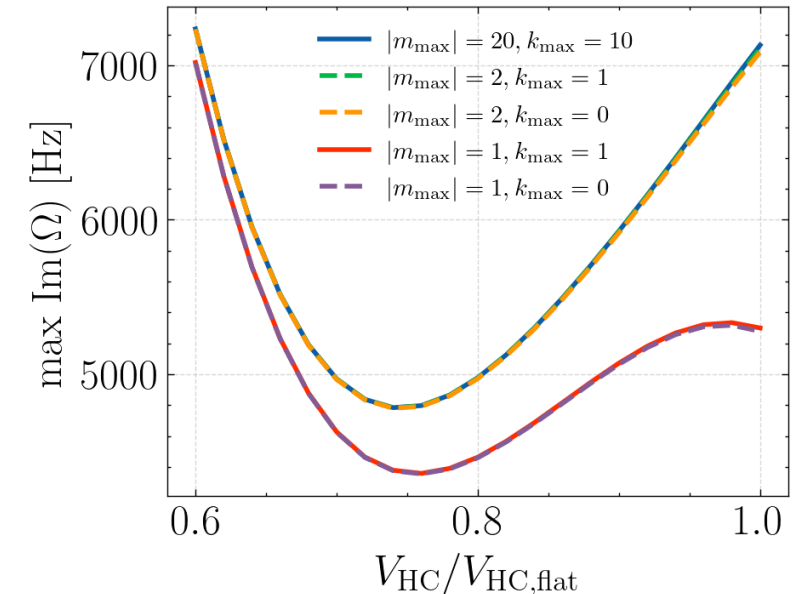
Only $|m|=1$ is insufficient
 $|m|=2$ required for instability

HALF, 350mA, SC HC



$|m|=1$ is sufficient
 $|m|=2$ required for accuracy

ALS-U, 500mA, NC HC



$|m|=1$ is sufficient for instability
 $|m|=2$ not very relevant

Maybe this explains why theories that consider only $|m|=1$ terms might fail in accurately predicting the mode 1 instability

Summary

- Tracking with one macroparticle predict mode 1 instability → simple mechanism
- Linear theory of mode-coupling with HC-modified parameters also predicts mode 1 instability. Good agreement with macroparticle tracking and experimental data.
- The mode 1 instability builds up when the coherent dipole frequency is too low (weak focusing)
 - Harmonic cavity → lower incoherent synchrotron frequency and longer bunch
 - Imaginary part of the impedance → negative shift to the coherent frequency
- This mechanism is insightful to understand the dependences of mode 1 instability/PTBL
 - Main RF cavity voltage
 - HC R/Q and detuning
 - Momentum compaction and energy spread
- A simplified theory that neglects non-linearity/Landau damping still provided useful information for the mode 1 instability with HCs. Results rely on the definition of a fictitious single-rf system that represents a double-rf system. Further studies to explore this approach are required.

Thank you for your attention

Acknowledgments

Fernando de Sá, Liu Lin, Ximenes Resende (LNLS)

Francis Cullinan, Åke Andersson (MAX-IV)

murilo.alves@lnls.br
cnpem.br



CNPq
Brazilian Center for Research
in Energy and Materials

MINISTRY OF
SCIENCE TECHNOLOGY
AND INNOVATION

